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# SiPM Readout Electronics

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## Abstract

Due to the peculiar characteristics of SiPM sensors in terms of equivalent capacitance, gain and fast rise-time response, in several applications classic readout solutions for radiation detectors are not able to provide optimal performance. Thus, several ad hoc readout approaches have been developed to fully exploit the favorable features of this kind of detectors. In this note the main requirements for the SiPM readout electronics are discussed, for both energy and time measurements, in the light of the detector model, and an overview of the main architectures commonly employed is provided, along with a set of relevant design examples.

*Keywords:* SiPM; front-end electronics; readout ASIC

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## 1. Introduction

Nowadays Silicon Photo-Multipliers (SiPM) represent a well consolidated and cost effective technology for a large range of applications requiring the detection of low light levels. In the last years, remarkable research efforts have been devoted, on the one hand, to improve the basic performance of this kind of detectors, for instance increasing the Photon Detection Efficiency (PDE), and, on the other hand, to reduce the impact of their

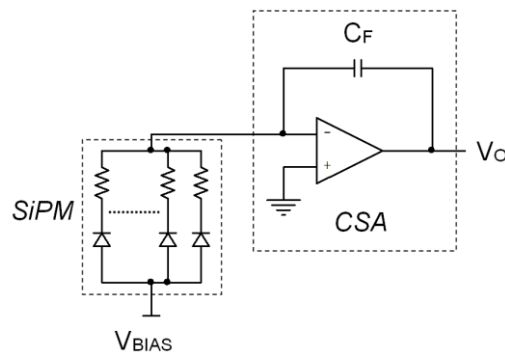
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18 main drawbacks, such as dark count rate, afterpulsing and optical crosstalk [1]. As a result, the possible  
19 application spectrum of SiPM detectors becomes wider and wider [2], covering fields where traditionally they  
20 have been considered a valid replacement for PMTs, such as Time of Flight Positron Emission Tomography  
21 (ToF-PET) [3] and calorimetry [4], but also more recently emerging areas, such as Light Detection and Ranging  
22 (LiDAR) [5] and Time Correlated Single Photon Counting (TCSPC) [6].

23 Often, front-end electronics plays a fundamental role in meeting the relevant specifications of a detection  
24 system based upon SiPMs and, in some cases, it even represents the bottleneck that limits the system  
25 performance. For instance, the Single Photon Time Resolution (SPTR) that is possible to achieve with a SiPM  
26 detector of large area is strongly dependent on the contribution of electronic noise [7], which causes statistic  
27 fluctuations of the instant when the output signal overcomes the chosen threshold. Thus, full exploitation of the  
28 favourable features of the detector requires the availability of suitable solutions for the front-end electronics,  
29 well-tuned to the peculiar characteristics of the SiPM. For energy measurements, often the main issue is not  
30 represented by the electronic noise, because of the large gain of the detector, around  $10^6$ . Typically, energy  
31 resolution is dominated by the intrinsic noise of the detector, associated to afterpulsing, optical crosstalk, dark  
32 pulses and gain fluctuations among the micro-cells. Moreover, in the applications where a SiPM is used to read-  
33 out a scintillator, the accuracy of the energy measurements is also strongly affected by the statistics of the  
34 photons released by the crystal.

35 The classic Charge Sensitive Amplifier (CSA) shown schematically in Fig. 1, which is the standard front-end  
36 configuration for radiation detectors, due to its very good noise performance, often does not represent an optimal  
37 solution for a SiPM.



38

39 Fig. 1. Charge Sensitive Amplifier (CSA) coupled to a SiPM.

40 In fact, several applications are characterized by large values of the maximum input charge, due to the high gain  
 41 of the detector, which makes direct charge integration on the feedback capacitance  $C_F$  impractical. Large values  
 42 of the integration capacitance are needed, to limit the charge-to-voltage gain  $1/C_F$  of the CSA, especially when  
 43 deep-submicron CMOS technologies are used and the allowed voltage headroom is very limited, due to the low  
 44 supply voltage. For instance, in a PET application with the detector coupled to an LSO crystal, the number of  
 45 photoelectrons contributing to the SiPM signal can reach 2000 [8] and, if the SiPM gain is  $10^6$ , this corresponds  
 46 to an input charge of 320pC. If the maximum output voltage swing of the amplifier in Fig. 1 is 1V, the feedback  
 47 capacitance needed would be 320pF, which is not feasible in a standard CMOS technology employed to design  
 48 multichannel readout ASICs. Even though the requirement in terms of maximum input charge is relaxed with  
 49 respect to the example above, the integration capacitance must be always quite large. This means that, if the  
 50 electronics must preserve the fast rising edge of the signal produced by the detector in order to achieve good  
 51 accuracy in time measurements, the amplifier must be able to drive large values of load capacitance while  
 52 featuring large bandwidth, which is possible only if power consumption is adequately increased. Other issues can  
 53 arise in terms of stability problems, once again due to large capacitive loads.

54 Concerning time measurements, in the vast majority of the SiPM applications Leading Edge Discrimination  
 55 (LED) is adopted as time pick-off method and the uncertainty in the evaluation of the occurrence time of the  
 56 detected event  $\sigma_t$  is related to the slope of the output signal of the front-end  $V_{OUT}$  around the chosen threshold  
 57  $V_{TH}$ , according to the well-known relation

$$58 \quad \sigma_t = \frac{\sigma_n}{\left. \frac{dV_{OUT}}{dt} \right|_{V_{OUT}=V_{TH}}} \quad (1)$$

59 where  $\sigma_n$  is the rms output noise of the front-end electronics [9]. The most relevant part of the charge released by  
 60 the SiPM is contained in the long tail of its current pulse, characterized by a very slow time constant, thus the  
 61 collection time of this fraction of the total charge is very long and the rise time of the output voltage of a CSA  
 62 coupled to the detector is determined by this slow time constant. As a consequence, a front-end solution based on  
 63 the CSA is not able to fully exploit the intrinsic fast leading edge of the current pulse generated by the SiPM and  
 64 requires the introduction of a fast shaper, which basically, starting from the output of the integrator, tries to  
 65 reconstruct a signal with fast rising edge, to be compared to a given threshold by means of a fast discriminator.

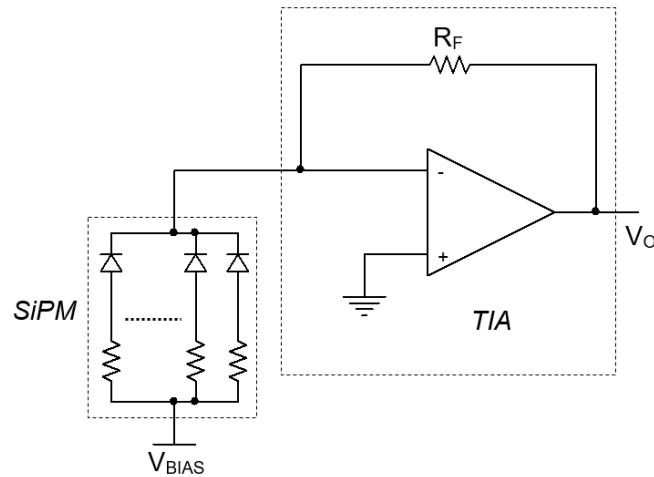
66 Instead of using a CSA, the most widespread approaches, adopted in several realizations of readout circuits for  
67 SiPM detectors, are based on interfacing the detector with a front-end at the same time able to preserve the  
68 intrinsic speed of the signal generated by the detector, thanks to well suited input impedance, and to reproduce at  
69 its output a replica of this signal, which can be conveniently applied to a fast discriminator for the extraction of  
70 the time information [10]. The choice of the most suitable solution for the electronics requires the availability of  
71 an accurate electric model of the SiPM, useful for reliably reproducing in simulations the interaction between the  
72 detector and the front-end electronics and much research work has been devoted in the past, and still is, to this  
73 task [11-13]. In any case a SiPM is always characterized by remarkable values of equivalent capacitance,  
74 especially when it is composed by a large number of micro-cells, thus the front-end architecture must cope with  
75 this feature.

76 The purpose of this note is discussing the issues related to the different approaches to the readout electronics  
77 for SiPMs, in the light of the characteristics of the detector, for both energy and time measurements. Some  
78 relevant solutions and realizations, representative of the different approaches, together with details about the  
79 implementation of the most relevant building blocks, will be also presented, trying to describe the progresses  
80 done and to understand the current trends.

## 81 **2. SiPM signal and front-end electronics**

82 As pointed out in the previous section, in many applications the classic CSA approach is not well suited for the  
83 readout of a SiPM, especially in multichannel integrated realizations using recent CMOS technologies, due to  
84 dynamic range, power consumption, stability and speed of response. Nevertheless, in applications characterized  
85 by limited input charge range, severe noise requirements and relaxed timing accuracy specifications, a CSA based  
86 front-end can be an interesting solution, due to its very good noise performance. Low gain SiPMs, with small  
87 micro-cell size and low total equivalent capacitance, can be conveniently read-out with a CSA preamplifier. An  
88 example of this kind of circuit is the ASIC VATA64HDR16 [14], which has been used, for instance, to read-out  
89 SiPM arrays coupled to continuous scintillation crystals for medical imaging applications [15] and in the  
90 detection of Cherenkov light, operating the detectors in photon counting mode [16].

91 Another possible solution, especially if timing accuracy is of interest, can be the transimpedance amplifier  
92 (TIA), depicted in Fig. 2.

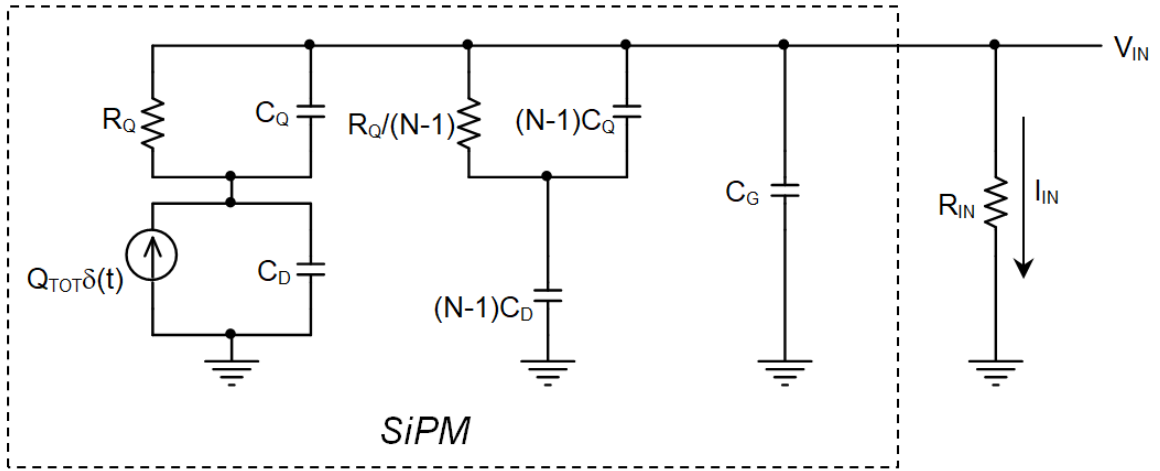


93

94 Fig. 2. Transimpedance amplifier (TIA) coupled to a SiPM.

95 The TIA converts the current pulse of the SiPM into a voltage and, if its bandwidth is large enough, it is able to  
96 preserve the fast rise time of the SiPM signal, thus enabling the achievement of good timing performance,  
97 according to (1). The main issue with this approach is once again the large equivalent capacitance of the detector  
98  $C_{DET}$ : if acceptable values of the current to voltage gain, i.e. of the feedback resistor  $R_F$ , are needed, the quite  
99 slow pole formed by  $R_F$  and  $C_{DET}$  at the input of the amplifier causes stability problems, which can be mitigated  
100 only by adding a compensation capacitance in parallel to  $R_F$ , thus limiting the closed loop bandwidth of the  
101 system, with detrimental effects on the timing accuracy. Moreover if also the energy information is needed, the  
102 output signal of the TIA must be integrated, which can be done with an integrator/shaper, often based on an  
103 active filter implementation that requires a voltage to current conversion. Thus, with this approach, the current  
104 pulse of the detector is first converted into a voltage, then back again into a current and finally integrated.

105 Let us now consider a classic model of the SiPM [11,12,17] coupled to a generic front-end preamplifier  
106 characterized by input impedance  $R_{IN}$ , depicted in Fig. 3.



107

108 Fig. 3. SiPM model coupled to a generic preamplifier with input impedance  $R_{IN}$ .

109 In Fig. 3,  $C_D$  is the capacitance of the photodiode in the micro-cell of the SiPM,  $R_Q$  is the quenching resistance,  
 110  $C_Q$  is the small parasitic capacitance in parallel to  $R_Q$  and  $C_G$  accounts for the total parasitic capacitance  
 111 associated to the large routing interconnections among all the micro-cells. In the left part of the figure a single  
 112 fired micro-cell is represented and the avalanche current is represented as an ideal Dirac's pulse containing the  
 113 total charge released by the detector

$$114 \quad Q_{TOT} = (C_D + C_Q)(V_{BIAS} - V_{BR}) = (C_D + C_Q)\Delta V \quad (2)$$

115 where  $V_{BIAS}$  is the bias voltage of the detector and  $V_{BR}$  is the breakdown voltage of the micro-cells. A Dirac's  
 116 delta can be conveniently used to describe the avalanche current, since avalanche breakdown is very fast as  
 117 compared to the time constant introduced by the circuit, especially considering the limitations inevitably  
 118 introduced by the bandwidth of the amplifier. In the central part of Fig. 3, the load of the other  $N-1$  micro-cells is  
 119 inserted. It is useful to notice that the same circuit can also be used to model the case in which more than one  
 120 micro-cell fires, by applying the superposition principle. The presence of  $C_Q$  accounts for the fast rising edge of  
 121 the voltage  $V_{IN}$  at the input of the amplifier, since, without this parasitic capacitance, the charge collection would  
 122 be dominated by the very slow time constant  $R_Q C_D$ . In other words,  $C_Q$  represents a fast path for the charge  
 123 generated by the avalanche towards the external circuit, since a fraction  $Q_F = Q_{TOT} C_Q / (C_D + C_Q)$  of the total charge  
 124 delivered by the avalanche flows almost instantaneously in  $C_Q$  and reaches very quickly the input node of the

125 preamplifier. An approximate analysis of the circuit provides the following expression of the contribution  $V_{INF}(t)$   
 126 of the “fast” charge  $Q_F$  to the input voltage of the preamplifier:

$$127 \quad V_{INF}(t) \cong \frac{Q_F}{C_{HF}} e^{-\frac{t}{\tau_F}}, \quad (3)$$

128 where, for low values of  $R_{IN}$ , which are commonly used in the applications as explained in the following, the time  
 129 constant  $\tau_F$  is

$$130 \quad \tau_F \cong R_{IN} C_{HF} \quad (4)$$

131 and  $C_{HF}$  is the equivalent capacitance of the detector for high frequencies

$$132 \quad C_{HF} \cong C_G + N \frac{C_D C_Q}{C_D + C_Q}. \quad (5)$$

133 In practice, eq. (3) states that the “fast” charge  $Q_F$  is almost immediately collected onto the high frequency  
 134 equivalent capacitance of the detector  $C_{HF}$ , which is then discharged on the input resistance of the preamplifier  
 135  $R_{IN}$ . Note that Eq. (3) works very well in the first few ns after the micro-cell is fired, since the capacitance  $C_{HF}$   
 136 can be conveniently used to describe the behavior of the detector only for fast transients. The input capacitance  
 137  $C_{IN}$  of the electronics does not play a relevant role, since normally it is negligible compared to  $C_{HF}$ .

138 Concerning the rest of the total charge  $Q_{TOT}$ , i.e.  $Q_D = Q_{TOT} C_D / (C_D + C_Q)$ , it is associated to the discharge  
 139 current  $I_D(t)$  of the capacitance  $C_D$  through the parallel connection of  $C_Q$  and  $R_Q$ , according to the slow recovery  
 140 time constant of the SiPM  $\tau_R = R_Q (C_D + C_Q)$ :

$$141 \quad I_D(t) = \frac{Q_D}{\tau_R} e^{-\frac{t}{\tau_R}}. \quad (6)$$

142 Considering again an approximate analysis of the circuit, the corresponding contribution to the input voltage of  
 143 the preamplifier is the following  $V_{INS}(t)$ :

$$144 \quad V_{INS}(t) \cong R_{IN} \frac{Q_D}{\tau_S - \tau_F} \left( e^{-\frac{t}{\tau_S}} - e^{-\frac{t}{\tau_F}} \right), \quad (7)$$

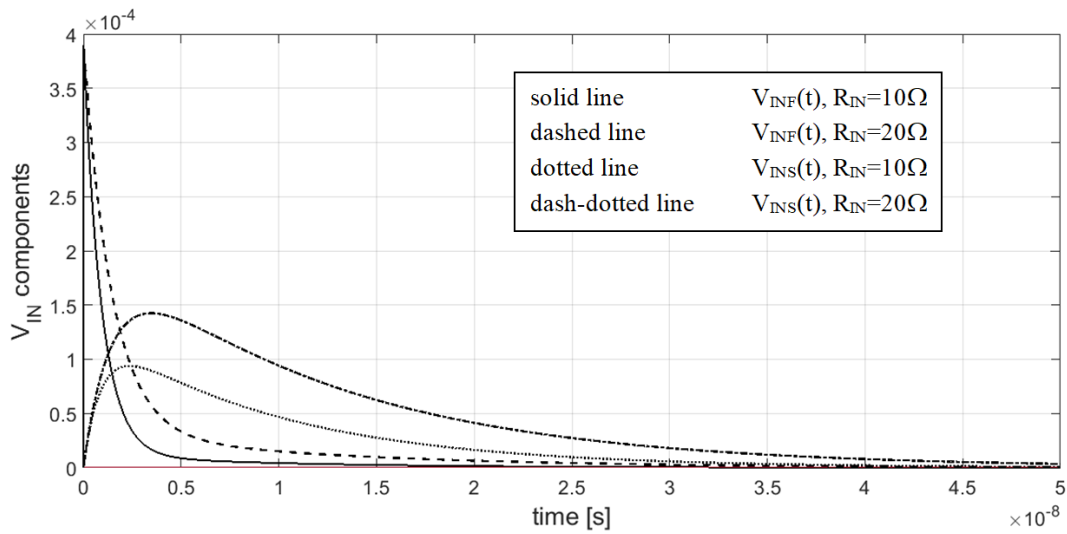
145 which exhibits a rise time dominated by the fast time constant  $\tau_F$  and a long tail dominated by the slow time  
 146 constant  $\tau_S$ :

$$147 \quad \tau_S \cong \tau_R + R_{IN} (C_G + N C_D) = \tau_R + R_{IN} C_{LF}, \quad (8)$$



148 being  $C_{LF}=(C_G+NC_D)$  the equivalent capacitance of the detector for low frequencies.

149 Of course, good timing performance can be achieved only thanks to the fast component  $V_{INF}(t)$  of the signal,  
 150 whereas the slow component  $V_{INS}(t)$  has poor relevance in this respect. Fig. 4 shows the results of a simulation of  
 151 the components of  $V_{INF}(t)$  and  $V_{INS}(t)$  for two different values of  $R_{IN}$ , i.e.  $10\Omega$  and  $20\Omega$ , considering values of the  
 152 parameters extracted for a SiPM with 3600 micro-cells [17], summarized in Table I.



153

154 Fig. 4. Simulation of the fast and slow components of the signal  $V_{IN}(t)$  for two different values of  $R_{IN}$ .

155 **Table 1**

156 SiPM parameters used for the SPICE simulations in Fig. 4

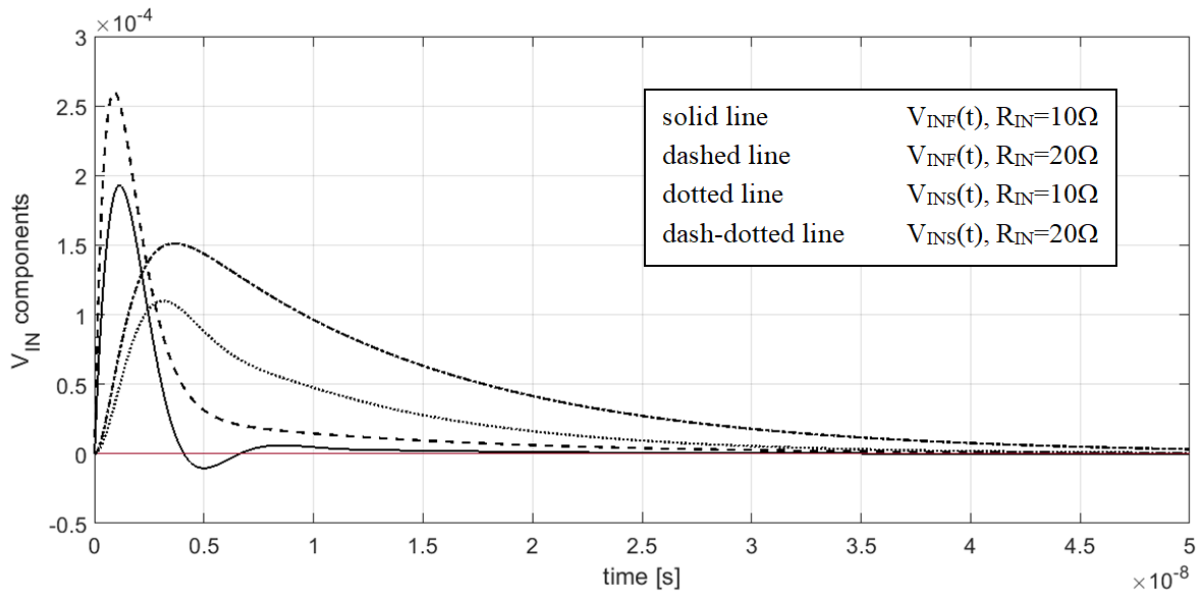
Parameter	Value
$C_D$	74fF
$C_Q$	30fF
$R_Q$	71k $\Omega$
$Q_{TOT}$	160fC

157

158 In Fig. 4 the effects of the variation of  $R_{IN}$  on the fast and slow components of the signal  $V_{IN}(t)$  are  
 159 highlighted. First, if  $R_{IN}$  is decreased, the contribution of the slow component  $V_{INS}(t)$  becomes less relevant and  
 160 the fall time of the fast component  $V_{INF}(t)$  gets faster. This means that the charge released by the detector is

161 collected more quickly if the input resistance of the front-end is reduced, because of the faster discharge of  $C_{HF}$ ,  
162 due to a larger discharge current flowing in  $R_{IN}$ . Moreover, the tail of the signal is apparently slower for larger  
163 values of  $R_{IN}$ . Consequently, if  $R_{IN}$  increases the rate of the event sustainable by the detection system is reduced,  
164 due to possible pile-up effects. In addition, the timing performance is affected, especially in case the time pick-off  
165 technique of choice is leading edge discrimination, since the time when the signal overcomes the threshold  
166 fluctuates, due to baseline variations. The previous considerations suggest that a very low input resistance is  
167 preferable for the front-end electronics of a SiPM detector. For instance, when the SiPM is used to read out a  
168 scintillator, low values of the input resistance are also required to limit the variation of the voltage across the  
169 detector as the photons impinge on it according to the characteristic time constant of the crystal. In fact, large  
170 voltage variations on the SiPM will cause non-linearity in the energy measurements, because the micro-cells  
171 undergoing avalanche breakdown at different times will experience appreciably different gain values.

172 The waveforms depicted in Fig. 4 are ideal, in that some important parasitics significantly affect, for instance,  
173 the rise time of the fast component  $V_{INF}(t)$ , limited only by the very fast time constants of the avalanche  
174 breakdown in the previous analysis. A remarkable contribution in this sense comes from the parasitic inductance  
175  $L$  associated to the interconnection between the detector and the front-end electronics [18, 19]. If the simulations  
176 in Fig. 4 are repeated in presence of an inductance  $L=10nH$ , the waveforms  $V_{INF}(t)$  and  $V_{INS}(t)$  are modified as  
177 shown in Fig. 5.



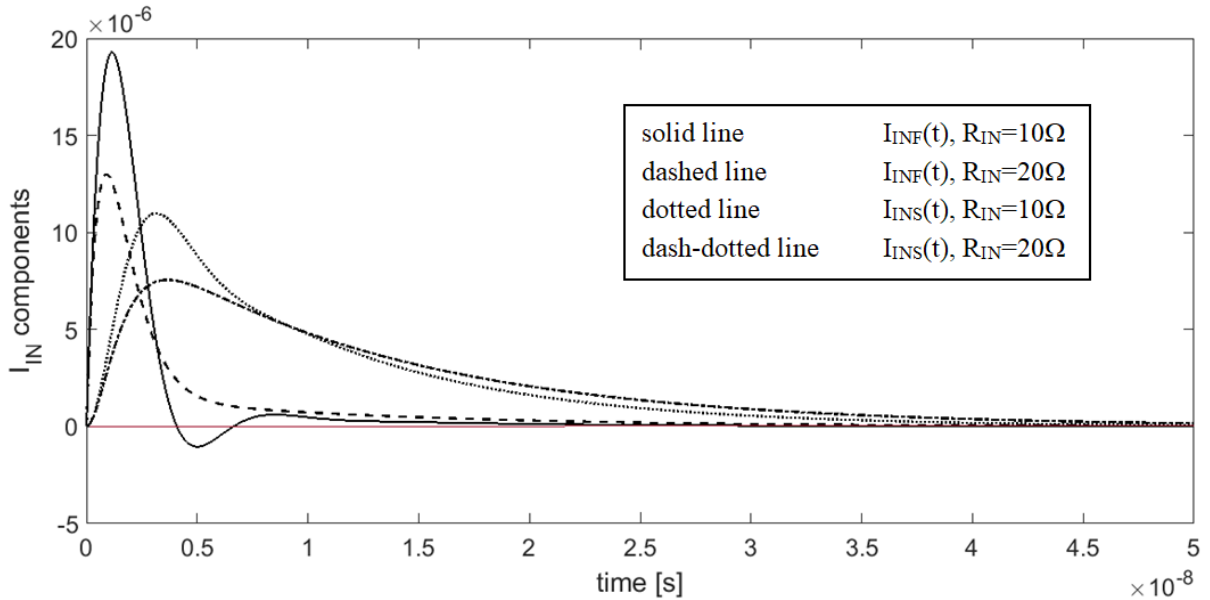
178

179 Fig. 5. Simulation of the fast and slow components of the signal  $V_{IN}(t)$  for two different values of  $R_{IN}$ , in presence of a parasitic  
180 interconnection inductance  $L=10\text{nH}$ .

181 On the one hand, the waveforms reported in Fig. 5 show that the slope of the fast component of the voltage  
182 signal is larger when  $R_{IN}$  is increased, which is good for timing accuracy. On the other hand, this feature cannot  
183 be fully exploited because once again increasing the resistance  $R_{IN}$  causes slower collection of the charge and  
184 longer signal tail also in presence of the parasitic  $L$ , as in the ideal case, leading to the same increased pile-up  
185 probability and non-linearity problems previously quoted. When the SiPM must be coupled to the electronics  
186 through a long cable, RF circuit techniques and  $50\Omega$  impedance matching can be applied to avoid signal  
187 reflections.

188 If the current  $I_{IN}$  which flows into the input resistance of the preamplifier is considered, some interesting  
189 conclusions can be drawn. The behavior of the fast and the slow components of this current,  $I_{INF}(t)$  and  $I_{INS}(t)$   
190 respectively, are shown in Fig. 6, obtained with simulations carried out in the same conditions of Fig. 5.  
191 Apparently, lower values of the input resistance correspond to both larger values of the slope of the fast  
192 component  $I_{INF}(t)$  and shorter tails for both components. This suggests that a very effective read-out approach can  
193 be based on a current mode preamplifier which reads the current pulse generated by the SiPM at very low  
194 impedance, discharging quickly the large equivalent capacitance of the detector  $C_{HF}$ , and reproduces this current

195 on a high impedance node, so that it can be further processed for the extraction of the time and energy  
 196 information. This is a very common approach for the front-end electronics used for SiPM and different  
 197 implementations of this scheme have been applied in several realizations of read-out circuits.

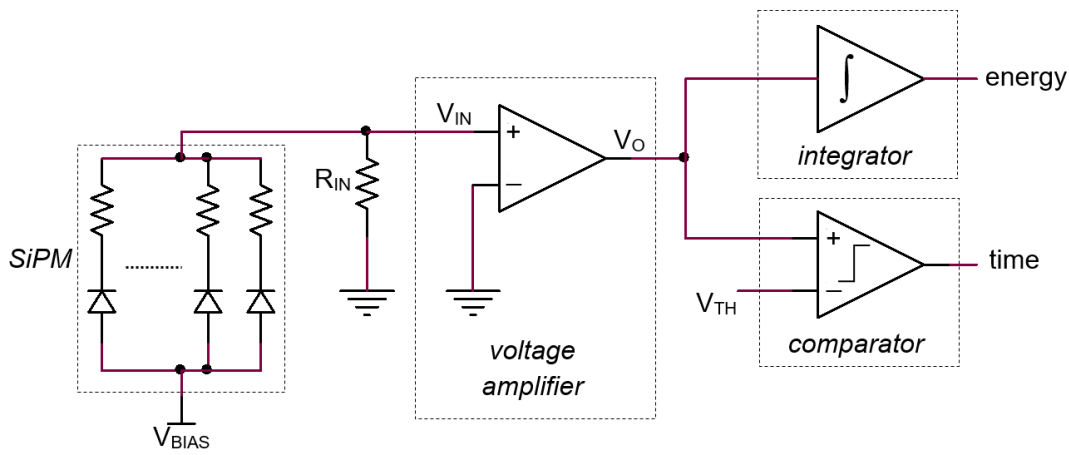


198  
 199 Fig. 6. Simulation of the fast and slow components of the current  $I_{IN}(t)$  for two different values of  $R_{IN}$ , in presence of a parasitic  
 200 interconnection inductance  $L=10nH$ .

201 **3. Voltage mode readout approach**

202 In any case, according to eq. (1) in order to achieve good timing performance, the preamplifier must be able to  
 203 preserve as much as possible the fast rise time of the input signal at its output, thus it must feature large  
 204 bandwidth. In case a voltage amplifier is used to read-out the detector, since  $R_{IN}$ , as discussed above, cannot be  
 205 large, the amplifier should also have sufficient gain, to reproduce an output signal of suitable amplitude, so that it  
 206 can be conveniently processed by the next blocks, e.g. an integrator and a comparator for energy or time  
 207 measurements respectively, as schematically depicted in Fig. 7. Such specifications, i.e. large gain-bandwidth  
 208 product, are difficult to be achieved without large power consumptions, making this approach not effective in  
 209 applications where very low levels of lights must be detected, timing accuracy is a relevant specification and the  
 210 number of readout channels is large. On the other hand, when the dynamic range of the input signal is large, thus

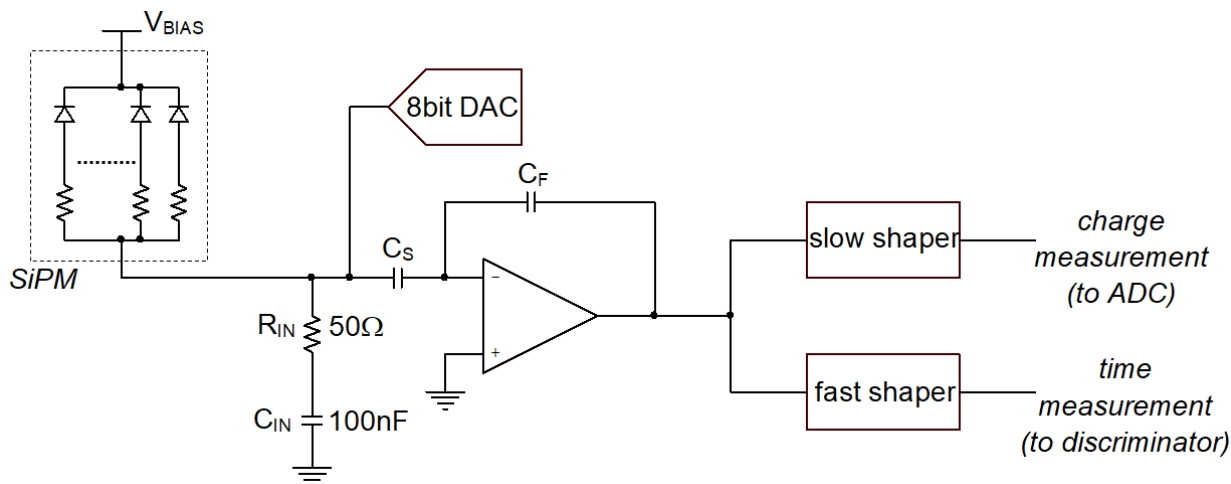
211 low voltage gain values are needed, and the specifications on the time accuracy are relaxed, the voltage mode  
 212 approach can be conveniently applied. As far as the noise performance of the circuit in Fig. 7 are concerned, the  
 213 total equivalent input voltage noise of the voltage amplifier, which is typically associated to a common source  
 214 input transistor, is directly summed to the voltage across  $R_{IN}$ , causing limitations in the Signal to Noise Ratio  
 215 (SNR) and in the timing resolution that is possible to achieve, according to (1).



216

217 Fig. 7. SiPM readout with voltage mode approach.

218 In the following, some examples of readout ASICs for SiPM detectors based on the voltage mode approach are  
 219 presented. SPIROC [20] and EASIROC [21], designed in  $0.35\mu\text{m}$  technology, exploit an external  $50\Omega$  input  
 220 resistance, DC decoupled, and an inverting voltage preamplifier to convert into a voltage and amplify the current  
 221 pulse of the SiPM. The preamplifier has the classic inverting feedback structure and the gain is set by means of  
 222 the capacitors  $C_S$  and  $C_F$ , as illustrated in Fig. 8 (SPIROC). Two different values of the gain can be chosen for the  
 223 preamplifier, i.e. 1 and 10.



224

225 Fig. 8. Structure of the analog channel of the ASIC SPIROC.

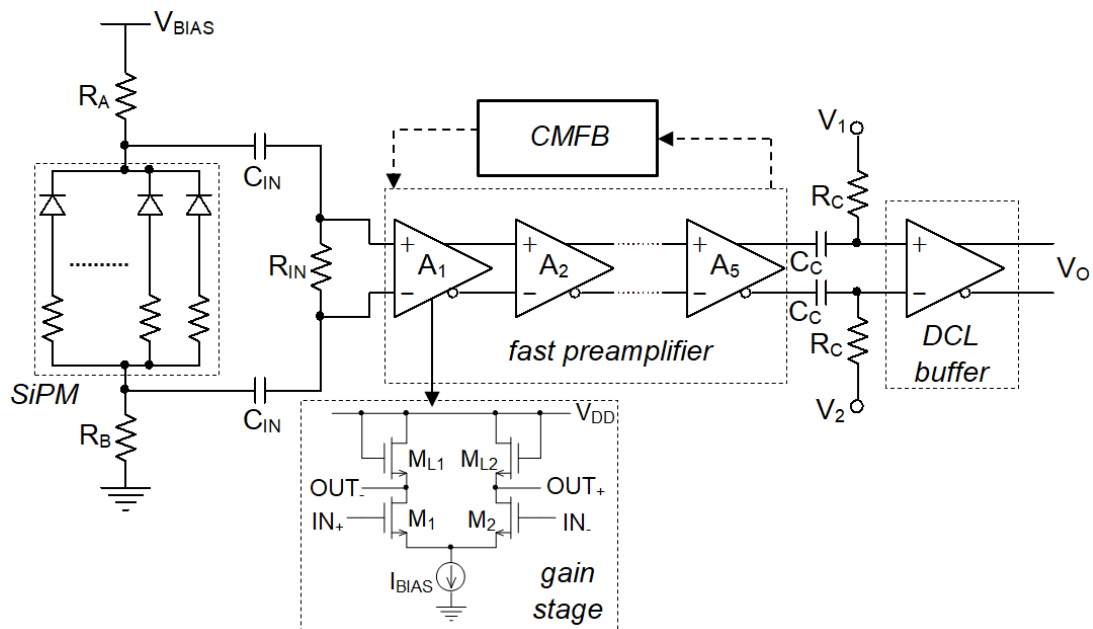
226 An 8 bit DAC is added at the input of the preamplifier to allow fine adjustments of the bias voltage of the  
 227 detector, which is DC coupled to the electronics. The CRRC<sup>2</sup> slow shaper integrates the output signal of the  
 228 voltage preamplifier with a selectable shaping time and, in parallel, a bipolar fast shaper, with shaping time of  
 229 15ns, is used to drive the fast voltage discriminator which generates the trigger signal.

230 The dynamic range of the charge measurements for both ASICs is 320fC, which corresponds to 2000 micro-  
 231 cells of a SiPM with gain 10<sup>6</sup>. The time resolution obtained is about 1ns rms for SPIROC, which exploits power-  
 232 pulsing techniques to reduce power consumption down to 25μW/channel, and is better for EASIROC, being  
 233 lower than 0.3ns rms for discriminator threshold set to 3pe (i.e. 0.48pC) and injected charges higher than 0.64pC,  
 234 with a power consumption of 4.84 mW/channel [22]. For both ASICs the threshold can be set down to 50fC,  
 235 making possible the detection of single photons. These circuits are not intended for applications with severe  
 236 requirements in terms of timing accuracy, such as ToF-PET, but are conveniently used in energy measurements  
 237 for physics experiments, for instance in calorimeters.

238 Considering the same ASIC family, in order to improve timing performance the PETIROC circuit [23] has  
 239 been designed in a SiGe technology, which offers HBT devices with  $f_t > 60\text{GHz}$  and allows achieving very large  
 240 bandwidth with limited power consumption. In this ASIC, in parallel with the inverting input stage followed by  
 241 the slow shaper used for the energy measurement, the signal path for the time measurements exploits an RF

242 common emitter preamplifier with a 10GHz gain-bandwidth product, followed by a fast voltage discriminator.  
 243 The measured jitter of the trigger signal generated with the front-end coupled to a 1x1mm<sup>2</sup> Hamamatsu MPPC,  
 244 using a fast laser to inject 15 photoelectrons and setting the threshold at 1photoelectron, is 46ps, with a power  
 245 consumption of 3.6mW/channel, excluding the output buffer used to observe the signal.

246 A further example of voltage mode preamplifier for SiPM detectors is the front-end of the first version of the  
 247 ASIC PETA [24] designed as an evolution of a previous circuit intended for photomultiplier tubes [25]. In this  
 248 circuit, the fast signal path used for time measurements is based on a fully differential voltage amplifier  
 249 composed by the cascade of 5 low-gain, high speed differential gain stages with diode loads, to maximize the  
 250 bandwidth [26] (see Fig. 9). The fully differential structure guarantees immunity from common mode noise and is  
 251 less sensitive to ground bounce and noise coming from the switching of digital parts, but more off-chip passive  
 252 components are needed (AC coupling) and the input pad number of the ASIC is doubled. The preamplifier reads  
 253 out the signal across an internal, adjustable termination resistance with nominal value of 50Ω. The gain of the  
 254 preamplifier is 20 and its maximum bandwidth is 900MHz, which can be limited in two steps by means of a low  
 255 pass filter. A slow common mode feedback block is also used to stabilize the common mode of the preamplifier  
 256 and a differential current mode logic (DCL) buffer, AC coupled to the preamplifier, acts as a fast discriminator.



257

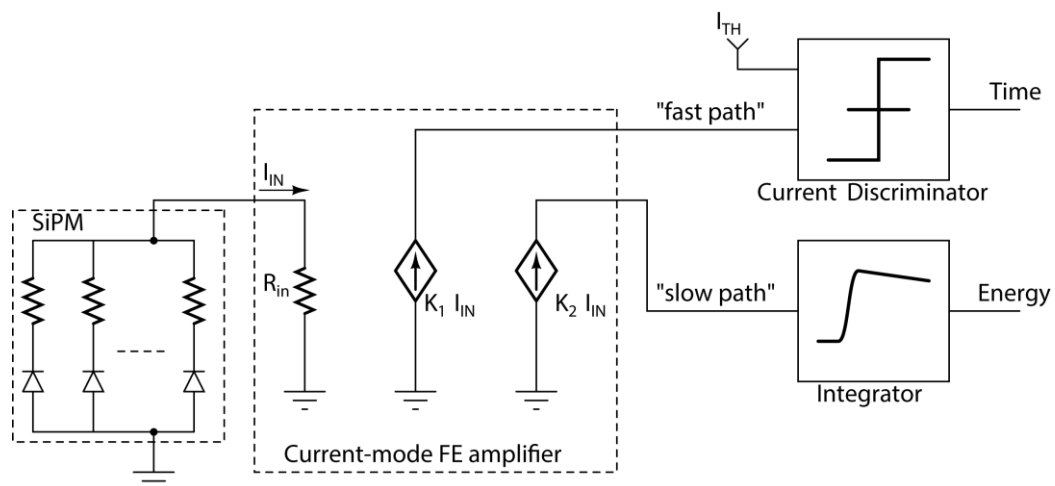
258 Fig. 9. Fast path of the ASIC PETA.

259 The PETA3 version of the ASIC provides very interesting results in terms of time resolution: the Coincidence  
 260 Resolving Time (CRT) obtained by coupling two channels of the ASIC to  $3 \times 3 \text{ mm}^2$  FBK SiPMs, used to read out  
 261  $3 \times 3 \times 5 \text{ mm}^3$  LYSO scintillators exposed to  $511 \text{ keV}$   $\gamma$ -rays, is  $190 \text{ ps}$  FWHM [27], but the power consumption is a  
 262 remarkable  $32 \text{ mW/channel}$ .

263 Always using a voltage mode approach, excellent timing accuracy can also be obtained with front-ends based  
 264 on commercial low-noise RF amplifiers intended for telecommunication and wireless applications, which exhibit  
 265  $50 \Omega$  input and output impedance and allow impedance matching, as mentioned in Section 2. For instance, in [28]  
 266 several experiments have been carried out with a monolithic RF amplifier with  $2 \text{ GHz}$  bandwidth, noise-figure of  
 267 about  $3.7 \text{ dB}$  at  $1 \text{ GHz}$  and gain of  $12 \text{ dB}$  (Minicircuits Mar-3SM+) coupled to a  $3 \times 3 \text{ mm}^2$  Hamamatsu MPPC and,  
 268 using  $2 \times 2 \times 3 \text{ mm}^3$  LSO scintillators, a very good CRT of about  $125 \text{ ps}$  has been achieved. However, also in this  
 269 case the power dissipation is huge, about  $400 \text{ mW}$ , and systems with a large number of channels are unpractical.

#### 270 4. Current mode readout approach

271 As pointed out in the conclusions of Section 2, current mode preamplifiers are commonly used to read out  
 272 SiPM detectors and also implementations with discrete components have been proposed in the literature [19]. Fig.  
 273 10 shows the basic principle of this approach: a current buffer with very low input impedance is coupled to the  
 274 detector and exploits the advantages of small  $R_{IN}$  values, illustrated in Section 2.



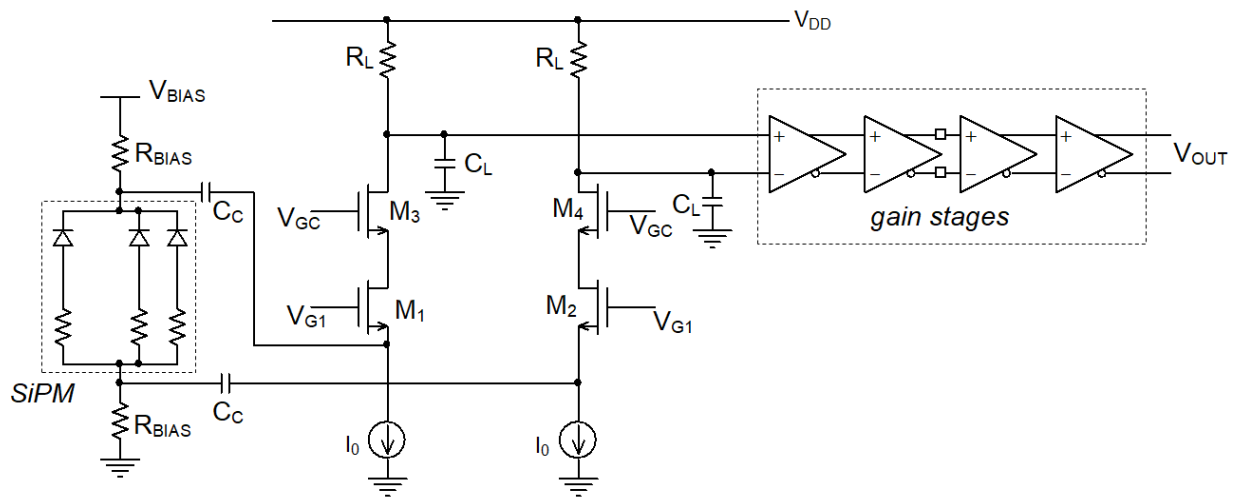
275

276 Fig. 10. Basic structure of a current mode analog channel for SiPM.



277 The output signal of the buffer is a high impedance replica of the current pulse generated by the detector that  
 278 can be easily reproduced with different scaling factors ( $K_1$  and  $K_2$  in Fig. 10) and used to establish different “fast”  
 279 and “slow” signal paths, optimized for charge or time measurements [29]. Typically, large bandwidths are easier  
 280 to be achieved with current mode amplifiers, because of the absence of high impedance nodes, thus the output  
 281 signal can follow the very fast leading edge of the current pulse generated by a SiPM, resulting in good  
 282 performance in terms of time resolution.

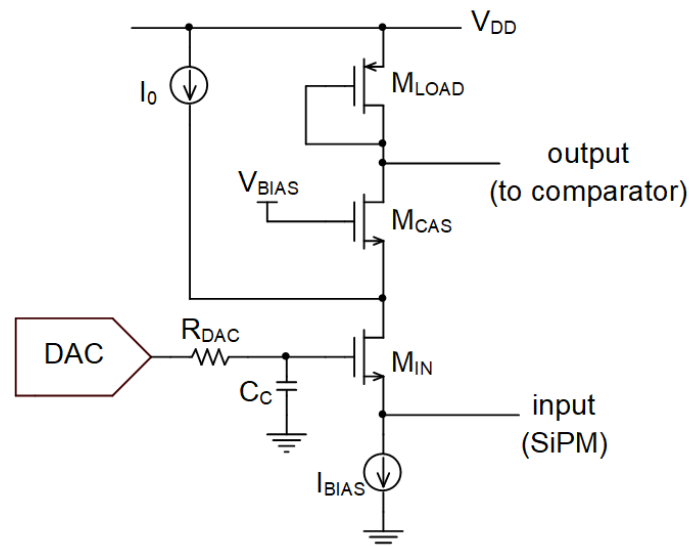
283 Several implementations of the current mode preamplifier have been proposed in the past. The simplest one is  
 284 based on a common gate current follower and one of the most relevant realization of this approach is the front-  
 285 end of the NINO ASIC [8,30], schematically depicted in Fig. 11.



286  
 287 Fig. 11. Analog channel of the ASIC NINO.

288 A fully differential configuration is employed to increase the immunity of the circuit against power supply and  
 289 ground noises. Cascode transistors  $M_3$  and  $M_4$  are used to decouple the drains of the input common gate  
 290 MOSFETs  $M_1$  and  $M_2$  from the output nodes and to increase the output resistance of the current buffer. The  
 291 current signal of the detector is converted into a voltage by means of the load resistors  $R_L$ , which form the  
 292 dominant time constant of the circuit (about 760ps) with the load capacitance  $C_L$  [31]. The input resistance of the  
 293 stage is set by the transconductance  $g_m=50\text{mA/V}$  of the common gate MOSFETs, corresponding to a total  
 294 differential resistance seen by the SiPM of  $40\Omega$ . The open loop configuration of the front-end makes the circuit

295 very fast and avoids any stability concerns. Considering a  $3 \times 3 \text{mm}^2$  Hamamatsu MPPC with an overvoltage of  
 296  $1.5 \text{V}$ , the rise time of the preamplifier output signal is  $1.5 \text{ns}$  for a single fired micro-cell and the estimated time  
 297 jitter, with a signal-to-noise ratio of 15, is  $100 \text{ps rms}$  [8]. The front-end is followed by four differential gain  
 298 stages, each with a voltage gain of 6 and  $500 \text{MHz}$  bandwidth, which form the discriminator, and the power  
 299 consumption of the preamp+discriminator is about  $20 \text{mW}$  [32]. Recent measurements carried out by coupling the  
 300 detector to a Hamamatsu S13360-3050CS SiPM biased at  $62 \text{V}$  achieve a Single Photon Time Resolution (SPTR)  
 301 of  $64 \text{ps rms}$  [33].  
 302 A similar configuration is adopted also by the circuit STiC3 [34], which also uses a differential front-end based  
 303 on an open-loop input common gate stage and a load resistor to form a voltage signal that is compared to a  
 304 threshold by means of a fast comparator for the generation of the trigger signal. In the last version of the circuit,  
 305 the load resistor has been implemented by means of a diode connected MOSFET [35] as depicted in Fig. 12 (only  
 306 one branch of the differential structure). The capacitance  $C_C$  allows keeping the value of the input resistance  $R_{IN}$   
 307 close to  $1/g_{mIN}$  also at high frequencies, mitigating the effects of the output resistance  $R_{DAC}$  of the DAC used to  
 308 fine tune the bias voltage of the SiPM. The results in terms of SPTR are very similar to the ones reported for the  
 309 NINO ASIC: using the same detector (Hamamatsu S13360-3050CS) an SPTR of  $67.1 \text{ps rms}$  has been achieved,  
 310 with a total power consumption of  $25 \text{mW/channel}$  [36].

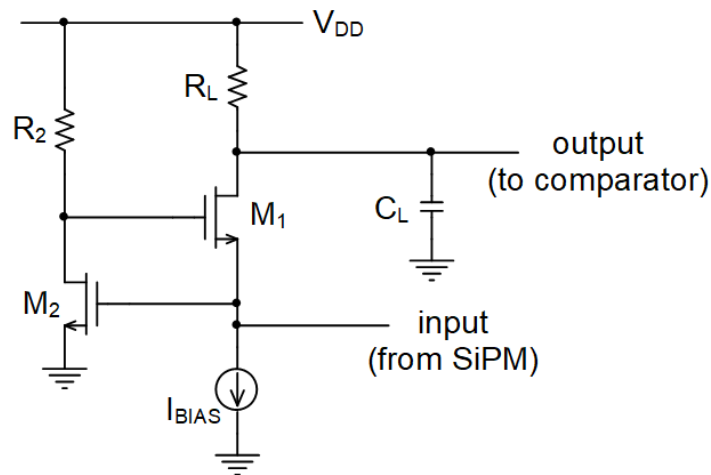


311

312 Fig. 12. Front-end of the ASIC STiC (only one branch of the differential current mode preamplifier).

313 Considering the noise performance of these preamplifiers, based on a common gate configuration biased by a  
314 current source, the contribution of the input transistor to the total noise at the output node of the preamplifier  
315 appears only at high frequencies, where the equivalent capacitance of the detector short-circuits to ground the  
316 source of the input MOSFET. At low frequencies, only the noise of the bias current source ( $I_0$  in Fig. 11 and  $I_{BIAS}$   
317 in Fig. 12) can reach the output node through the common gate transistor, but this contribution can be made  
318 negligible by reducing the transconductance of the MOSFET which provides the bias current. As a result, in  
319 general the common gate stage exhibits better noise performance as compared to a common source preamplifier,  
320 very often used in the voltage mode readout approach. This advantage of the current mode readout in terms of  
321 noise, which translates in better timing performance due to eq. (1), decreases when large detectors, characterized  
322 by large equivalent capacitance, are used. On the other hand, the output dynamic range of the common gate  
323 amplifier is reduced by the presence of the bias current source, whereas this limitation does not apply to a  
324 common source stage, thus, in principle, we can conclude that the current mode approach is worse than the  
325 voltage mode in terms of dynamic range.

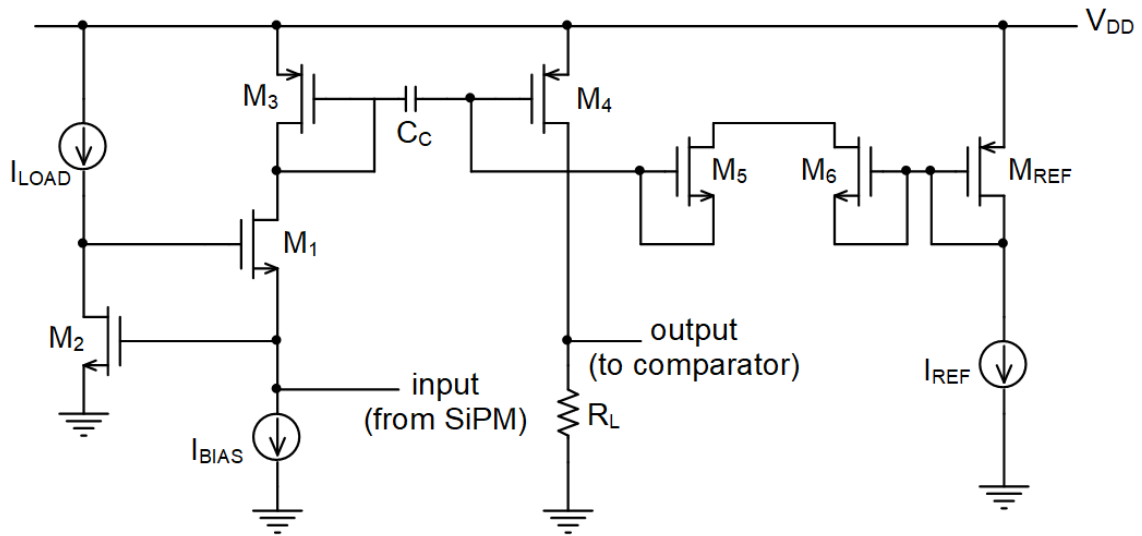
326 Feedback is often employed in order to decrease the input resistance of the current mode preamplifier while  
327 saving power. One of the most used topologies is the regulated common gate transimpedance amplifier,  
328 schematically depicted in Fig. 13. Here, the open-loop input resistance of the preamplifier  $1/g_{m1}$  is decreased by a  
329 factor equal to the loop gain of the applied feedback  $g_{m2}R_2$ , thus there are more degrees of freedom, with respect  
330 to the simple common gate solution, that can be exploited to find a good compromise solution among low  $R_{IN}$ ,  
331 power consumption and noise. Also stability is of concern in this kind of circuit and one of the main design  
332 guidelines is limiting the Q factor of the complex conjugate poles of the circuit to values less than 1, to control  
333 the damped oscillatory behavior of its impulse response. Detailed circuit analysis can be found in [37] and shows  
334 that the output noise is dominated by the contribution of the noise current of  $M_2$ , which can be reduced by  
335 increasing  $g_{m2}$  with respect to  $g_{m1}$ , thus ascribing to the feedback the task of reducing the input resistance of the  
336 preamplifier.



337

338 Fig. 13. Regulated common gate as a front-end for SiPM.

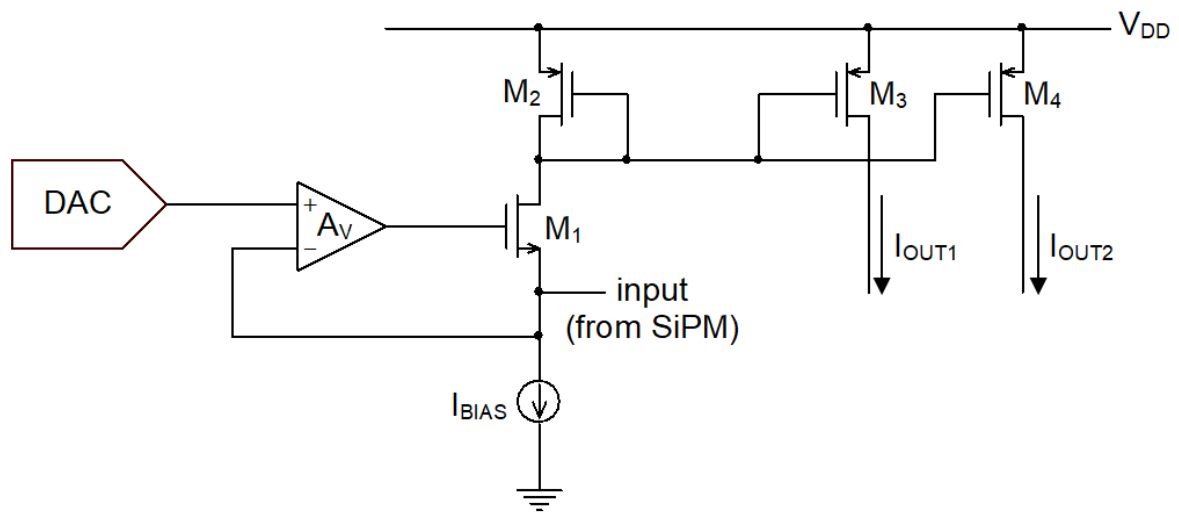
339 An example of read out ASIC intended for SiPM detectors and based on the regulated common gate approach  
 340 is the TOFPET2 circuit [38]. The front-end can be coupled to both n-on-p and p-on-n devices: Fig. 14 shows the  
 341 basic circuit used for one of the two SiPM polarities. The load of the regulated common gate is the diode  
 342 connected PMOSFET  $M_3$ , which is AC coupled to a common source amplifier with passive load, to convert the  
 343 current pulse of the detector into the input voltage of the discriminator. The operating point of the common  
 344 source  $M_4$  is settled by means of the voltage reference formed by  $I_{REF}$  and  $M_{REF}$ , connected to the gate of  $M_4$  via a  
 345 large resistor, in the order of  $G\Omega$ , realized by the back-to-back cut-off MOSFETS  $M_5$  and  $M_6$ . The same  
 346 arrangement, basically consisting in an AC coupled current mirror, is replicated to obtain different signal paths  
 347 with suitable scaling factors, to be used for energy and time measurements. As an example of the timing  
 348 performance of the TOFPET2 ASIC, an SPTR of 95 ps rms has been measured using a fast laser source and one  
 349 of the SiPM of the Hamamatsu 4x4 array HPK S13361-3050AE-04, biased at 7.5V overvoltage [39]. The overall  
 350 power consumption of the ASIC is less than 10mW per channel.



351

352 Fig. 14. The regulated common gate preamplifier of the ASIC TOFPET2.

353 A slightly different implementation of the regulated common gate that can be found in several realizations is  
 354 the one depicted in Fig. 15 [40-43].



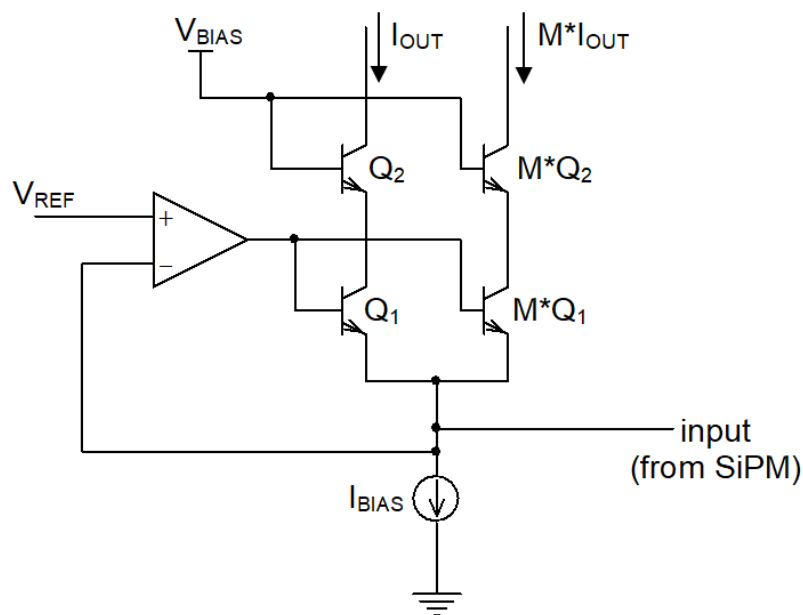
355

356 Fig. 15. Regulated common gate front-end with gain boosting realized by means of a differential amplifier.

357 Here the input resistance of the preamplifier is given by  $R_{IN}=1/g_{m1}A_V$ , where  $A_V$  is the gain of the differential  
 358 gain boosting stage in the regulation feedback loop. Also in this case the stability issue must be carefully taken  
 359 into account in the design and the gain-bandwidth product of the differential amplifier must be large, to guarantee

360 low values of the input impedance of the preamplifier also at high frequencies. The main advantage associated to  
361 the structure of Fig. 15 is the possibility of fine tuning the bias voltage of the SiPM by exploiting the virtual short  
362 circuit at the terminals of the differential amplifier, using a DAC.

363 As already mentioned, in several cases two different signal paths are established and optimized to carry out in  
364 effective way both time and energy measurements. In the architectures considered up to now, these paths are  
365 formed after the front-end stage, for instance exploiting current mirrors. In [44] a different approach is proposed:  
366 the current pulse of the SiPM is split into two fractions directly at the input of the preamplifier, by exploiting two  
367 matched HBTs, available in the SiGe technology used to design the circuit, suitable scaled of a factor  $M$  and  
368 arranged in a regulated common base configuration, as shown in Fig. 16. The same configuration has been also  
369 used for the input stage in [45], using MOSFETs instead of HBTs.



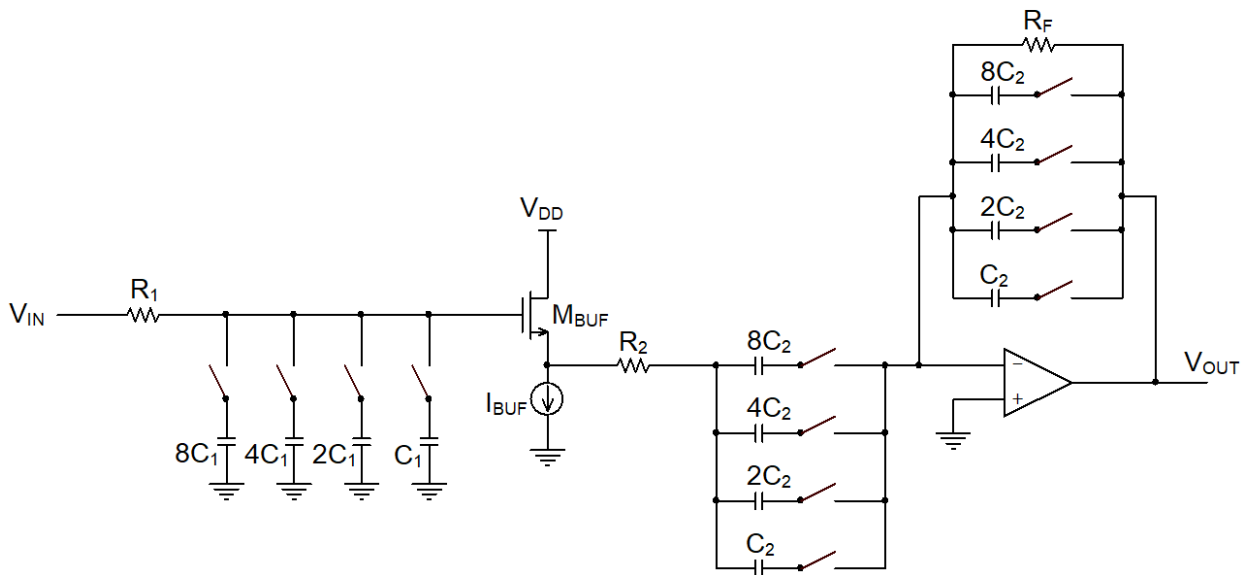
370  
371 Fig. 16. Two signal paths formed directly at the input of the preamplifier by means of matched HBTs arranged in a regulated common base  
372 configuration.

373 Several further examples of current mode front-ends for SiPM that exploit feedback in different ways to  
374 decrease the input resistance of the circuit can be found in the literature [45-50], thus we can conclude that the

375 current mode approach is the most common solution used for the readout of this kind of detectors, especially in  
376 applications where timing resolution is of interest, for instance ToF-PET.

### 377 5. Energy measurements: main circuit solutions

378 As discussed in the previous sections, the very front-end of an electronic channel intended for SiPMs must  
379 preserve as much as possible the favourable features of the detector and provides an output signal proportional to  
380 the current pulse of the detector. In order to extract the information about the charge associated to the signal,  
381 proportional to the energy of the detected event, the most straightforward approach is integration of the signal  
382 itself. In case a voltage mode preamplifier has been used, integration can be performed by means of a slow shaper  
383 cascaded to the front-end. The shaper can be implemented with a passive RC network or, more frequently, with  
384 an active filter. For instance, in the ROC ASIC family, a CRRC<sup>2</sup> shaper has been cascaded to the inverting  
385 voltage preamplifier in the slow path of the signal. The shaping time is variable between 25ns and 200ns, to  
386 accommodate the requirements of different applications [51]. Fig. 17 shows an example of the structure of the  
387 slow shaper with programmable shaping time [52].



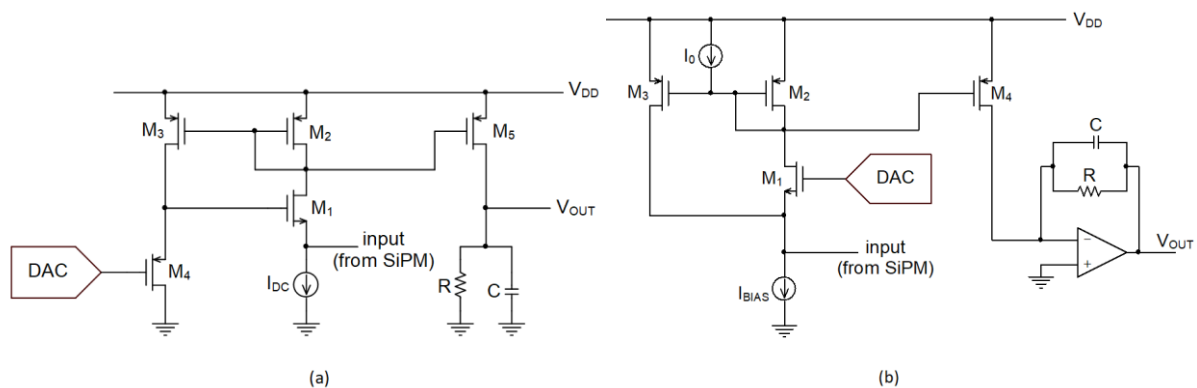
388

389 Fig. 17. Configuration of the CRRC<sup>2</sup> slow shaper used to integrate the output signal of the SPIROC preamplifier.

390 In the circuit of Fig. 17, the first pole is realized by means of a passive low-pass RC circuit and the rest of the  
 391 transfer function of the shaper has been implemented using an active band-pass filter, decoupled from the passive  
 392 RC network by means of a voltage buffer.

393 For instance, in PET applications the signal in response to a 511keV event is the result of the convolution  
 394 between the single-photon response of the SiPM, characterized by its long tail, and the function of time which  
 395 describes the photon emission rate of the scintillator, characterized by its time constant. The peaking time should  
 396 be long enough to integrate as much as possible the resulting signal, for accurate evaluation of the energy, and  
 397 typical values in this case are around 200ns, for LYSO or LSO scintillators. Of course long shaping times limit  
 398 the event rate that can be sustained by the channel.

399 In case a current mode approach is used for the front-end, integration of the current signal of the fast path, as  
 400 in Fig. 9, is a straightforward task and can be carried out by means of a passive RC network [45,47] or using a  
 401 CSA [29]. These two solutions are illustrated in Fig. 18, which represents the front-end and the integrator of the  
 402 ASIC Klaus (Fig. 18a) [47] and BASIC (Fig. 18b) [29]. In Fig. 18 the different approaches applied in these two  
 403 cases for the application of feedback to the input current buffer can be distinguished.

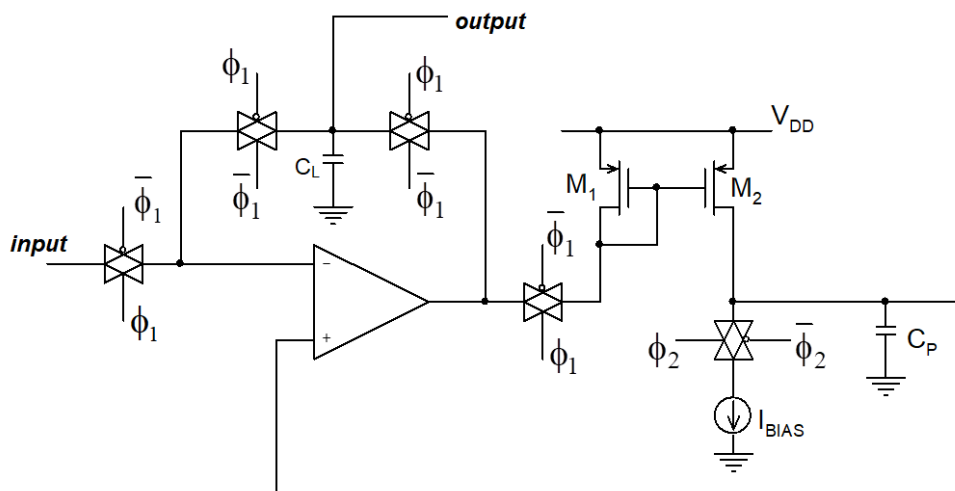


404 (a) (b)  
 405 Fig. 18. Front-end and slow path of the analog channels of the ASICs Klaus (a) and BASIC (b).

406 In both voltage and current mode, the peak of the integrator output is the signature of the energy to be  
 407 measured, thus it must be sampled and made available for further processing, for instance analog to digital  
 408 conversion. This can be accomplished by using an external signal to sample the integrator output voltage at the  
 409 peaking time of the shaper, generated by means of a delay after the fast discriminator has fired. This sampling

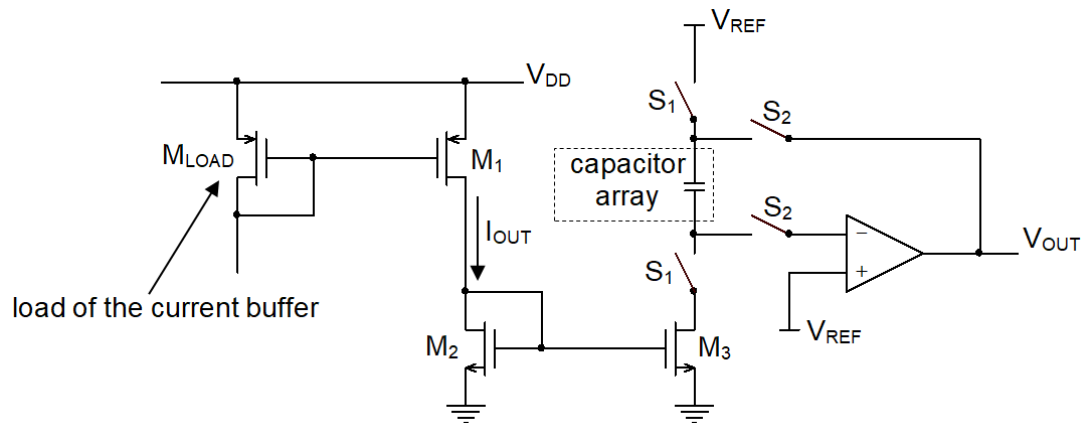


410 signal can be used to store the integrator output in an analog memory realized with a switched capacitor array  
 411 properly addressed, as in the ROC family [20,21,23]. As an alternative, a peak stretcher circuit can be used to  
 412 detect and store the peak voltage of the integrator [29,48]. This circuit is often realized according to the structure  
 413 proposed in [53] and schematically represented in Fig. 19. The current mirror  $M_1$ - $M_2$  works as a rectifying  
 414 element in a feedback loop and the circuit can be easily configured so that, while waiting for a valid signal, the  
 415 voltage on the capacitor  $C_P$  follows the input voltage ( $\phi_1=1$ ;  $\phi_2=1$ ). As soon as a valid event has been detected, the  
 416 current  $I_{BIAS}$  is switched off ( $\phi_1=1$ ;  $\phi_2=0$ ) and the  $C_P$  cannot be discharged, thus the voltage across  $C_P$  tracks the  
 417 peak of the input voltage. Last, when the peak is reached, the OPAMP goes into positive saturation, the current  
 418 mirror is switched off and the circuit is reconfigured as an analog memory ( $\phi_1=0$ ;  $\phi_2=0$ ), presenting at the output  
 419 the voltage across  $C_P$ .



420  
 421 Fig. 19. Peak detection circuit.

422 Another approach consists in integrating the output signal of the front-end in a time window of suitable  
 423 duration, often programmable, started when the fast discriminator fires [24,38,54]. In the ASIC TOFPET2 [38],  
 424 which features a front-end based on a current mode approach as described in Section 4, the output current of the  
 425 front-end is switched on a capacitor array, via current mirrors, to integrate the current pulse when the fast  
 426 discriminator fires ( $S_1$  closed,  $S_2$  open in Fig. 20).

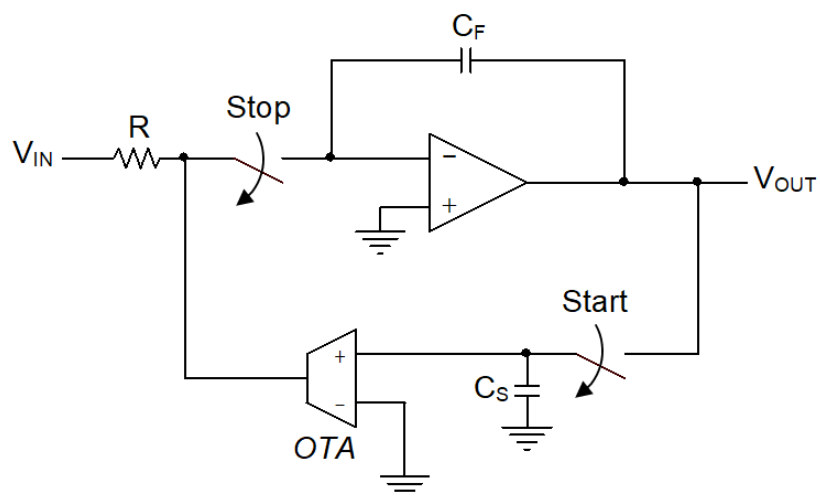


427

428 Fig. 20. Slow path of the ASIC TOFPET2.

429 At the expiration of the integration time window, the capacitor array is switched on a differential output buffer  
 430 ( $S_1$  open,  $S_2$  closed).

431 Instead, in [24] the front-end output is a voltage pulse, which is converted into a current and integrated in the  
 432 chosen time window using a classic CSA without resistive feedback, as illustrated in Fig. 21. A slow feedback  
 433 loop around the integrator, based on an Operational Transconductance Amplifier (OTA), is used to compensate  
 434 the DC component of the input voltage. This loop is opened as soon as the Start signal, provided by the fast  
 435 discriminator, is activated, so that the output current of the OTA stays constant and only the AC component of the  
 436 input voltage is integrated. The integration time window is closed by the Stop signal, generated by a timer circuit.

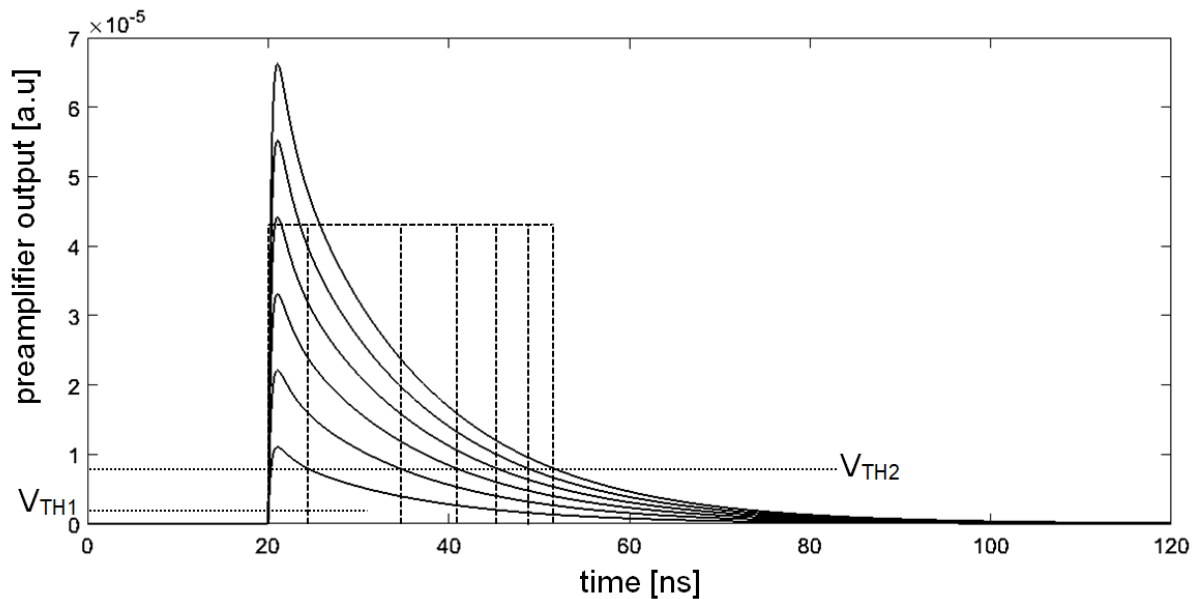


437

438 Fig. 21. Integrator of the ASIC PETA.

439 Often, the integrator is followed by an ADC and the energy information is made available in digital format  
440 [20,38,55], so that it is easily transferred to the external electronics for further processing via fast digital links.  
441 All the approaches discussed so far guarantee very good linearity in energy measurements, which is needed in  
442 applications such as gamma spectroscopy or calorimetry.

443 When linearity is not of particular concern, one of the most common techniques for evaluating the energy of  
444 the detected events is Time over Threshold (ToT), which consists in measuring the duration of the preamplifier  
445 output pulse, associated to the charge generated by the detector. Usually, the pulse duration is evaluated as the  
446 time interval in which the pulse amplitude is comprised between two different thresholds, as depicted in principle  
447 in Fig. 22: the first one,  $V_{TH1}$ , is the same very low threshold of the fast discriminator, used for the measurement  
448 of the event occurrence time, whereas the second one,  $V_{TH2}$ , is higher, in order to improve the energy resolution  
449 which can be achieved.

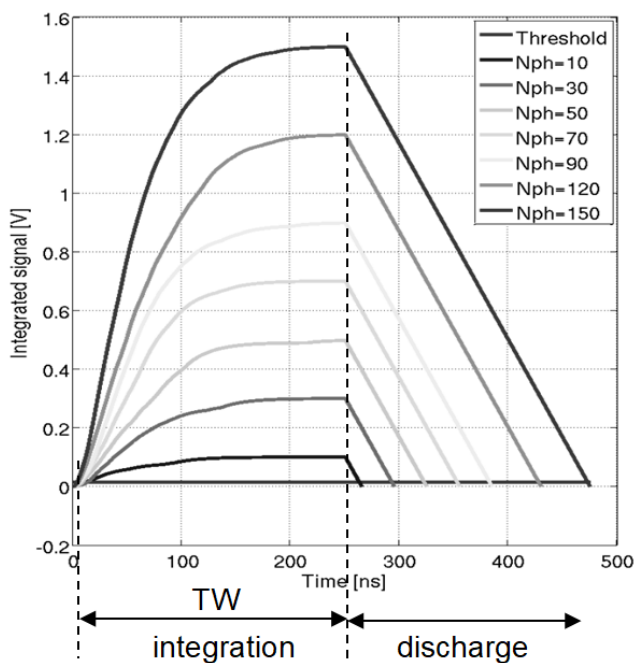


450  
451 Fig. 22. Time over Threshold (ToT) technique for energy measurements.

452 In this way, also the energy measurements are carried out by means of time measurements and, for instance, a  
453 Time to Digital Converter (TDC) can be exploited for both time and energy evaluation, resulting in more  
454 compact electronics and power consumption saving. In any case the relationship between the duration of the  
455 pulse obtained with this technique and the charge generated by the detector is strongly non-linear, as Fig. 22

456 clearly shows. For instance, in [8] a differential passive filter has been interposed between the SiPM terminals  
 457 and the input of the NINO preamplifier, in order to shape the waveform of the current pulse, so as to mitigate  
 458 non-linearity of the ToT technique and extend the dynamic range of the charge that can be processed. In general,  
 459 ToT can be conveniently exploited in applications that do not require much accuracy in the energy measurements,  
 460 in order to optimize the requirements of the electronics in terms of compactness and power consumption. For  
 461 instance, this technique is applied in PET applications, for identification of the photo-peak and correction of the  
 462 time-walk in time measurements [32,42].

463 To improve the linearity of the ToT technique for charge measurements, another approach can be exploited:  
 464 the output pulse of the front-end, converted into a current in case a voltage mode preamplifier is used, is  
 465 integrated within a time window of suitable duration TW, using a capacitor. Then, the integration capacitor is  
 466 discharged with a constant current down to the threshold, obtaining a ramp signal with duration proportional to  
 467 the total integrated charge. This technique, illustrated in Fig. 23, has been applied, for instance, in [46,56,57] and  
 468 makes possible the exploitation of a TDC for energy measurements also in applications based on a continuous  
 469 crystal read-out by an array of SiPMs, in which the energy resolution achieved with a simple ToT solution is not  
 470 sufficient, due to the large statistic dispersion of the long tail of the SiPM signal.



471

472 Fig. 23. Linearized Time over Threshold (ToT) technique for energy measurements: fixed integration time  $T_W$  and variable discharge time  
473 for signals corresponding to increasing number of photoelectrons  $N_{ph}$ .

474 In [35] a similar technique exploits the saturation of the input common gate stage, which turns off when the  
475 amplitude of the detector current increases, to integrate the charge on the same capacitance of the detector. Then  
476 this capacitance is discharged by the bias current of the input stage.

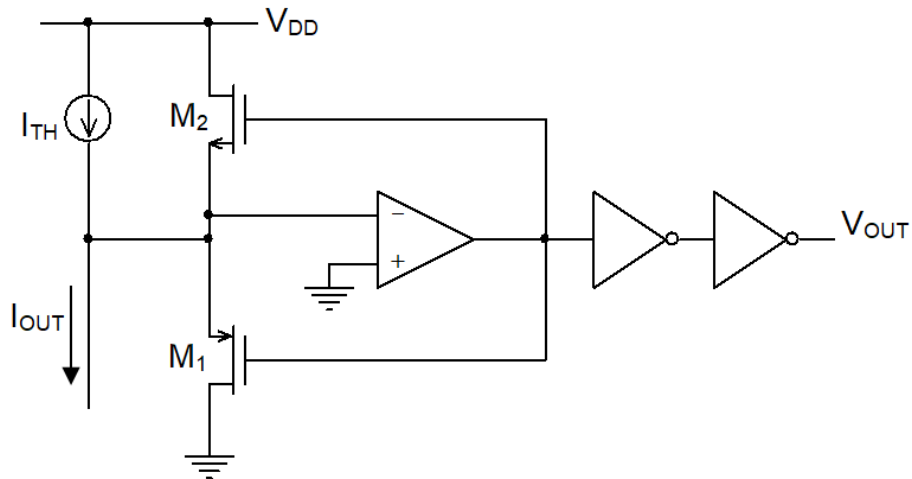
477 As a general observation concerning the accuracy of energy measurements with the ToT, the slope of the slow  
478 falling edge of the SiPM pulse is very small, especially when the number of photons to be detected is low.  
479 According to eq. (1), this affects the time jitter of the output signal of the discriminator and makes problematic  
480 the recourse to the ToT techniques when the application requires the identification of the number of photons and  
481 single photon spectrum is needed.

## 482 **6. Time pick off solutions**

483 As already pointed out in the introduction, LED is the favourite time pick-off method with SiPMs and in  
484 several read out circuits, based on both voltage and current mode, a fast voltage comparator is used to form an  
485 output signal with a very sharp transition when the detector pulse overcomes the threshold. This trigger signal  
486 marks the arrival of the event and can be time-stamped by means of a Time to Digital Converter (TDC). In some  
487 cases [8,24], as already shown, the discriminator is composed by the cascade of low gain, large bandwidth  
488 voltage amplifiers, with overall gain sufficient to generate a fast, full swing output pulse in response to the signal  
489 generated by a single micro-cell of the SiPM undergoing avalanche breakdown. Hysteresis can be added to the  
490 discriminator by means of a small amount of positive feedback, to avoid undesired output transitions due to the  
491 noise [8].

492 In current mode front-end circuits, a current discriminator is often used in the fast signal path to compare the  
493 output current pulse of the front-end to a threshold current, thus no further current to voltage conversion is  
494 needed, as in case a voltage discriminator is employed. A common structure for the fast discriminator is the one  
495 proposed in [58], schematically represented in Fig. 24. In DC, the PMOS  $M_1$  is ON and carries the difference  
496 between the threshold  $I_{TH}$  and the output current of the front-end  $I_{OUT}$ , so that the output voltage of the inverting  
497 amplifier is low. As soon as  $I_{OUT}$  overcomes the threshold, due to the arrival of a valid event,  $M_1$  turns off and the

498 NMOS  $M_2$  turns on, thus the output of the amplifier makes a fast positive transition, sensed by the cascaded  
 499 inverters. Examples of readout circuits which exploit this current comparator structure are, for instance,  
 500 [29,40,49,59,60].



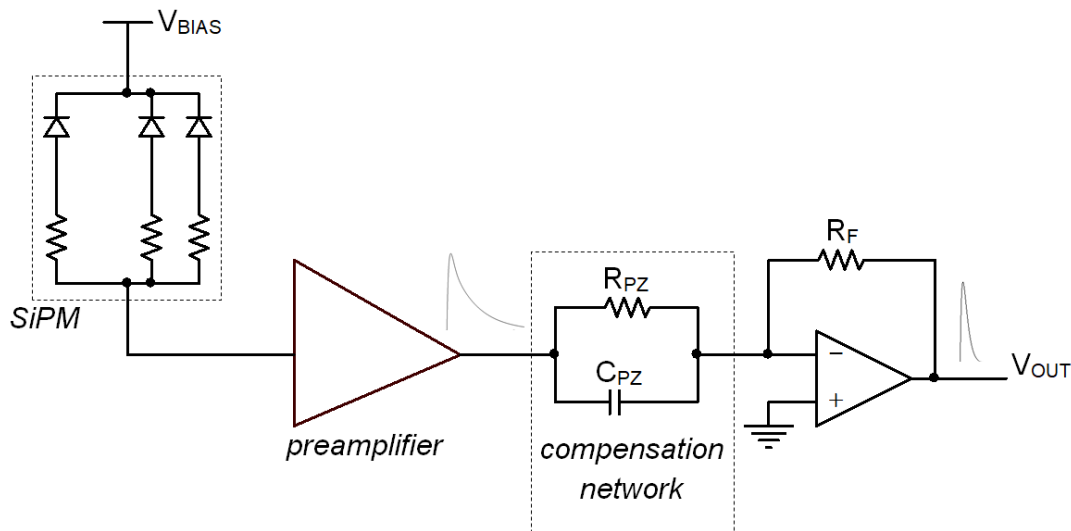
501

502 Fig. 24. Structure of a current discriminator.

503 Time walk, i.e. the dependence of the time when the threshold is overcome on the amplitude of the signal, is a  
 504 typical issue of the LED time pick-off technique, and the classic circuit solution for this problem is Constant  
 505 Fraction Discrimination (CFD) [9]. However, very few examples of application of this technique to SiPM readout  
 506 exist. For instance, in [61] a comparison is reported between the results obtained in terms of timing accuracy by  
 507 using LED and CFD applied to the output of the same current mode front-end. When the front-end is coupled to a  
 508  $3 \times 3 \text{ mm}^2$  Hamamatsu MPPC, used to readout a  $3 \times 3 \times 15 \text{ mm}^3$  LFS crystal, a time resolution of 479ps and 712ps  
 509 FWHM have been achieved with LED and CFD respectively. Non-linearity of the current differentiation stage  
 510 used in the CFD circuit and large variability of the rise time of the SiPM signal are identified as the main causes  
 511 of the worse results obtained with CFD. Moreover, in relevant applications, such as PET, only the events around  
 512 the photo-peak are of interest, thus the amplitude variability is not the main cause of errors in time measurements.

513 Another interesting method proposed to improve time resolution in SiPM readout systems used for PET is the  
 514 Differential Leading Edge Discrimination (DLED) technique [62]. This method is useful to get rid of the baseline  
 515 fluctuations due to the long tail of SiPM dark pulses, which causes errors in the evaluation of the time when a

516 valid event overcomes the discriminator threshold. The tail of the SiPM signal, characterized by the slow time  
 517 constant  $\tau_s$ , is compensated by means of a linear filter introduced in the fast signal path, which adds a zero with  
 518 the same time constant to the overall transfer function of the front-end. A simple example of this pole-zero  
 519 cancellation filter is reported in Fig. 25. This results in a fast return to zero of the dark pulses generated by the  
 520 detector, without affecting the rise time of the response to the single photon. A downside of this method is that  
 521 the differentiation introduced by the filter tends to increase the contribution of the electronic noise in eq. (1), thus  
 522 the DLED technique cannot be used when this effect becomes relevant. Moreover the time constant of the zero  
 523 must be adjustable if the readout circuit is intended to be coupled to different kinds of SiPM.



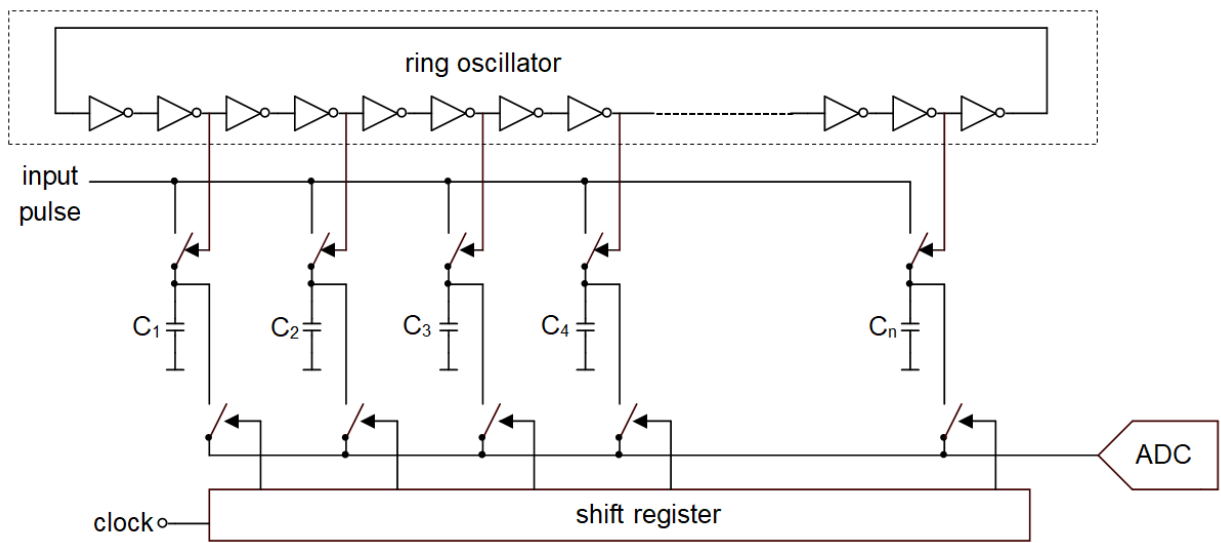
524

525 Fig. 25. The DLED technique, based on the compensation of the SiPM tail by means of pole-zero cancellation.

526 A similar technique is presented in [63], where the compensation zero is introduced by choosing a suitable  
 527 value for the capacitor used for AC coupling of the detector to the front-end electronics, according to the value of  
 528 the input resistance  $R_{IN}$  of the front-end.

529 Very good time resolution can be obtained with more sophisticated time pick-off techniques, based on digital  
 530 processing of a number of samples of the SiPM pulse [64,65]. Several examples of multichannel fast digitizer  
 531 ASICs have been proposed in the literature [66-69]. The basic circuit structure exploited in this class of circuits is  
 532 a fast Switched Capacitor Array (SCA), and a simple example of architecture is illustrated in Fig. 26 [70]. The  
 533 sampling signal propagates through the inverter chain, which forms a ring oscillator, and the capacitors are used

534 as analog memories to store the samples of the signal. The depth of the capacitor array allows the storage of the  
 535 SiPM pulse and, after the sampling phase, the shift register allows the readout of the array towards an ADC,  
 536 which can be internal or external. As an example of performance of this kind of circuits, the DRS4 ASIC [69]  
 537 hosts 8+1 channels composed by an array of 1024 storage cells; the sampling rate can be varied from 700MS/s  
 538 and 6GS/s, the analog bandwidth is 950MHz and the power consumption is between 10 and 40mW/channel,  
 539 depending on the sampling speed and the selected mode of operation.



540

541 Fig. 26. An example of fast sampler: Domino structure.

542 Concerning the digital time pick-off methods that can be used if several samples of the raising edge of the  
 543 detector pulse are available, there is a broad range of solutions. A large class of the techniques are an extension of  
 544 the corresponding analog ones and exploit the samples of the signal in order to reduce the effects of noise and  
 545 make more accurate the evaluation of the time when the threshold is crossed. For instance, in Digital Constant  
 546 Fraction Discrimination the time when the signal overcomes a fixed fraction of the pulse amplitude is computed  
 547 using the measured samples of the pulse [71]. In [72] the two consecutive samples of the pulse raising edge  
 548 characterized by the maximum difference (maximum slope) are found and the time of the event is obtained by  
 549 means of the intersection of the line which connects these samples with the baseline of the pulse.



550 Interpolation of the available samples around the threshold and normalization of the pulse amplitude is often  
551 used to increase the total number of samples and improve the accuracy in the determination of the threshold  
552 crossing time: for instance in [73] cubic spline interpolation is applied for this task.

553 Another class of methods is based on true digital algorithms. A possible approach tries to find a matching  
554 between a reference pulse, evaluated by means of real data or by theoretical analysis, and the measured samples,  
555 in order to reconstruct the start time of the event from the start time of the matched reference signal. For instance  
556 in [74], the pulse is modelled with the sum of two exponentials and the least mean square difference between the  
557 normalized measured pulse and the reference pulse progressively shifted is minimized. Deconvolution by means  
558 of an optimal filter is also used to reduce the effects of noise [73,75]. The application of digital processing  
559 techniques requires remarkable computing resources, depending on the nature and the complexity of the chosen  
560 technique, thus much effort is devoted in trying to simplify the algorithms employed and to make them suitable to  
561 be implemented on compact and easy to use devices, such as FPGA and DSP.

## 562 **7. Conclusions**

563 A review of the main approaches commonly applied for the readout electronics dedicated to SiPM detector has  
564 been proposed and the related issues have been discussed. A comparison between voltage mode and current mode  
565 front-end circuits has been done in the light of a simple model of the detector, showing why the latter is now the  
566 favourite choice, especially in the integrated realizations. The most common solutions used for energy and time  
567 measurements have been also presented and compared. Concerning the development perspectives in the field of  
568 front-end electronics for SiPMs, accurate modelling of the detector is one of the main issues that is still open and  
569 the role of the parasitic elements, associated to the interconnections between the SiPM and the electronics, in the  
570 formation of the signal deserves more detailed studies, so that an optimal choice for the specifications of the  
571 preamplifier, such as its input resistance and bandwidth, can be made. Another relevant issue is the evolution of  
572 the integrated technologies towards nanometer devices, mainly oriented to digital applications and very  
573 challenging when analog circuits have to be designed. For instance power supply reduction is one of the main  
574 concerns and require the development of suitable solutions to preserve the slope of the output pulse of the  
575 preamplifier, thus the accuracy in time measurements and, at the same time, a large dynamic range for the charge

576 measurements. On the other hand the resort to this kind of technologies opens opportunities in terms of possible  
577 integration of more digital resources on chip, making possible the realization of very compact Systems on Chip  
578 (SoC) and increasing the data communication bandwidth. In this respect, a very important factor that must be  
579 taken into account is the cost of the development of electronics in these nanometer technologies. The required  
580 investments are increasing more and more and accessibility to this kind of technologies will be probably limited  
581 only to large collaborations.

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